

# Molecular Vibrations

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## Why Estimate Molecular Vibrations?

- Simulation of vibrational spectrum (identification of molecules)
- Vibrational corrections to enthalpy
- (Small) vibrational corrections to polarizability and other properties
- Understanding of vibrational motion could assist dynamics experiments and “mode-selective” chemistry

## Small Vibrations in Classical Mechanics

The classic reference is Wilson, Decius, and Cross, *Molecular Vibrations* (Dover, New York, 1980). Cheap book, makes a good reference.

Let us focus on purely classical systems at first; all the results carry over to quantum mechanics.

For small vibrations, the motion of atom  $\alpha$  away from its equilibrium value may be described by  $\Delta x_\alpha, \Delta y_\alpha, \Delta z_\alpha$ , with kinetic energy

$$T = \frac{1}{2} \sum_{\alpha=1}^N M_\alpha \left[ \left( \frac{d\Delta x_\alpha}{dt} \right)^2 + \left( \frac{d\Delta y_\alpha}{dt} \right)^2 + \left( \frac{d\Delta z_\alpha}{dt} \right)^2 \right]$$

If we switch to mass-weighted coordinates, such as  $q_1 = \sqrt{M_1}\Delta x_1$ ,  $q_2 = \sqrt{M_1}\Delta y_1$ ,  $q_3 = \sqrt{M_1}\Delta z_1$ ,  $q_4 = \sqrt{M_2}\Delta x_2$ , etc., then the kinetic energy operator becomes simpler since the mass factors are now absorbed

$$\begin{aligned}
 T &= \frac{1}{2} \sum_{i=1}^{3N} \dot{q}_i^2 \\
 V &= V_0 + \sum_{i=1}^{3N} \left( \frac{\partial V}{\partial q_i} \right)_0 q_i + \frac{1}{2} \sum_{i=1}^{3N} \left( \frac{\partial^2 V}{\partial q_i \partial q_j} \right)_0 q_i q_j + \dots
 \end{aligned} \tag{1}$$

Remember that at equilibrium,  $(\partial V/\partial q_i)_0 = 0$ ; we can also set  $V_0 = 0$ . Also abbreviate  $(\partial^2 V/\partial q_i \partial q_j)_0$  as just  $f_{ij}$ .

## Newton's Equations of Motion

We can rewrite Newton's equations of motion as

$$\frac{d}{dt} \frac{\partial T}{\partial \dot{q}_j} + \frac{\partial V}{\partial q_j} = 0 \quad j = 1, 2, \dots, 3N$$

or

$$\ddot{q}_j + \sum_{i=1}^{3N} f_{ij} q_i = 0$$

A possible solution to this equation is

$$q_i = a_i \cos(\sqrt{\lambda} t + \phi)$$

where the angular frequency is  $\sqrt{\lambda}$ ; this is just  $\sqrt{k/m}$  in harmonic oscillator — the  $m$  has been absorbed by the mass-weighted coordinate system used here!

Substitute the last expression into the differential equations to get

$$\sum_{i=1}^{3N} (f_{ij} - \delta_{ij}\lambda) a_i = 0 \quad j = 1, 2, \dots, 3N$$

or in matrix notation, just  $\underline{\underline{\mathbf{F}}} \underline{\mathbf{a}} = \lambda \underline{\mathbf{a}}$ . *This is an eigenvalue equation!* We have a solution to this system of  $3N$  linear equations only if  $\lambda$  has special values obtainable from the secular determinant

$$\begin{vmatrix} f_{11} - \lambda & f_{12} & f_{13} & \cdots & f_{1,3N} \\ f_{21} & f_{22} - \lambda & f_{23} & \cdots & f_{2,3N} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ f_{3N,1} & f_{3N,2} & f_{3N,3} & \cdots & f_{3N,3N} - \lambda \end{vmatrix} = 0$$

## Normal Modes of Vibration

The matrix eigenvalue equation is equivalent to matrix diagonalization which is equivalent to solving the secular determinant for each  $\lambda$  ( $N$  of them). Once we have the eigenvalues  $\lambda_k$  we can get the corresponding eigenvectors  $\underline{\mathbf{a}}_k$ , giving the motion of each atom for the given eigenvalue  $\lambda_k$ :

$$q_{ik} = a_{ik} \cos \left( \sqrt{\lambda_k} t + \phi_k \right).$$

The eigenvectors  $\underline{\mathbf{a}}_k$  are the *normal modes* of vibration. *For each normal mode, all the atoms move with the same frequency and phase, but with different amplitudes.*

## Normal Coordinates

We can define a new set of coordinates using the normal modes. This gives us the “normal coordinates”

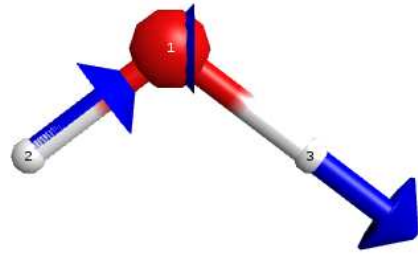
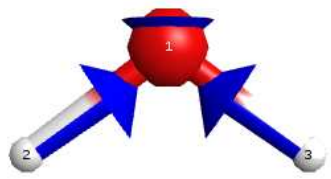
$$Q_k = \sum_{i=1}^{3N} a_{ik} q_i \quad k = 1, 2, \dots, 3N$$

Since the eigenvectors of a real, symmetric matrix (**F**) are orthogonal,  $T$  and  $V$  become diagonal (no cross terms):

$$T = \frac{1}{2} \sum_{k=1}^{3N} \dot{Q}_k^2$$
$$V = \frac{1}{2} \sum_{k=1}^{3N} \lambda_k Q_k^2$$

The Hamiltonian is separable in this representation!





## Polyatomic Molecules

What happens for quantum mechanics, and for polyatomic molecules? Use Harmonic Oscillator model.

- $3N-6$  frequencies ( $3N-5$  for linear molecules); the rest are translations and rotations with zero frequency
- In normal mode coordinates, Hamiltonian is separable: wavefunction is a product and energy is a sum. Total vibrational energy is  $\sum_i \omega_i \hbar (v_i + 1/2)$
- Minimum energy (due to uncertainty principle) is “zero point vibrational energy” (ZPVE or ZPE), where  $v_i = 0$  for all  $i$ .  $ZPVE = \frac{1}{2} \hbar \sum_i \omega_i$

## How Would We Get Harmonic Frequencies for a Molecule?

- Easy — just diagonalize the second derivative matrix  $\underline{\underline{\mathbf{F}}}$ , called the *Hessian*. The frequencies  $\omega_i$  are the square roots of the eigenvalues,  $\sqrt{\lambda_i}$ .
- Where do we get  $\underline{\underline{\mathbf{F}}}$ ?
- Recall  $f_{ij} = (\partial^2 V / \partial q_i \partial q_j)$
- Potential energy  $V$  is just  $E_e$  (B.O. approximation!): Need  $\partial E_e^2 / \partial q_i \partial q_j$ .
- Compute second derivative of  $E_e$  in terms of Cartesian displacements ( $x_\alpha, y_\alpha, z_\alpha$ , call them  $\tilde{q}_i$ ) and it's easy to transform to mass-weighted coordinates, using

$$\underline{\underline{\mathbf{F}}} = \underline{\underline{\mathbf{M}}}^{-1/2} \underline{\underline{\tilde{\mathbf{F}}}} \underline{\underline{\mathbf{M}}}^{-1/2}.$$

- How do we get  $\partial^2 E_e / \partial x_\alpha \partial y_\beta$ , etc? Need second derivative of electronic energy vs nuclear coordinates. Can compute *analytically* (using formula) or *numerically* from finite differences of energies or gradients:

$$\frac{\partial^2 E_e}{\partial x_\alpha \partial y_\beta} \approx \left[ \left( \frac{\partial E_e}{\partial y_\beta} \right) \Big|_{x_\alpha = x_{\alpha 0} + \Delta x_\alpha} - \left( \frac{\partial E_e}{\partial y_\beta} \right) \Big|_{x_\alpha = x_{\alpha 0} - \Delta x_\alpha} \right] / \Delta x_\alpha.$$

## Analytic Hessian Better than Numerical

- Analytic Hessian might cost  $\sim 10\text{-}30\text{x}$  cost of energy; analytic gradient costs maybe  $\sim 1.5\text{-}2\text{x}$ .
- Can need gradients from many displaced geometries — up to  $6N$  (+ and - for each of  $3N$  coordinates) — unless reduced by point group symmetry
- Numerical Hessian contains numerical errors (divide small number by small number)
- (Can land on wrong solution if displacement drops symmetry)

## Availability of Analytic Derivatives

Method	Gradient	Hessian
HF, DFT	Y	Y
CI	Y	N
CCSD, CCSD(T)	Y	S
MP2	Y	S
CASSCF	Y	S
CIS	Y	Y
EOM-CCSD	S	N
TD-DFT	S	N

S = available in some packages; Y = widely available

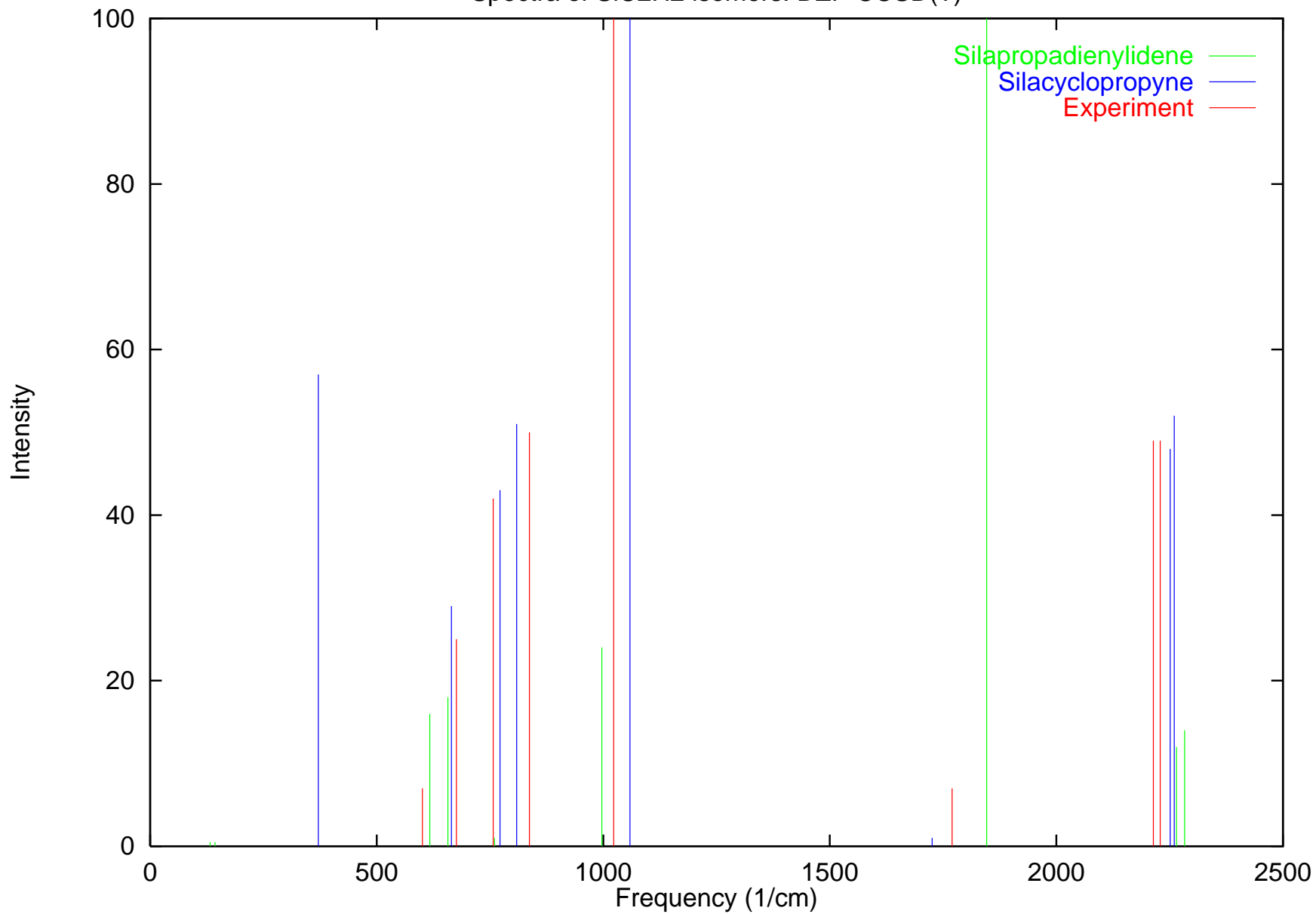
## Approximate Average Errors in Harmonic Frequencies

(Using polarized double and triple zeta basis sets)

Method	Error
HF	11%
CISD	4-6%
CCSD	1-4%
CCSD(T)	1-3%

Anharmonicity accounts for another  $\sim 2-3\%$  difference from experimental fundamental frequencies. Many workers employ *scaling factors* for each level of theory to better predict fundamental frequencies.

Spectra of SiC<sub>2</sub>H<sub>2</sub> isomers: DZP CCSD(T)





## Scaling ZPVE's

In an enlightening paper, Grev, Janssen, and Schaefer [*J. Chem. Phys.* **95**, 5128 (1991)] showed that using scaled fundamental frequencies to estimate the ZPVE is *not* necessarily better than using unscaled frequencies. The reason is anharmonicity.

If ZPVE's use scaling, they should have a *different* scaling factor than the individual frequencies.

$$G(v) = \sum_r \omega_r \left( v_r + \frac{1}{2} \right) + \sum_{r \geq s} \chi_{rs} \left( v_r + \frac{1}{2} \right) \left( v_s + \frac{1}{2} \right) + \dots,$$

$$\Delta^{harm} = G(0) - ZPVE^{harm} = \frac{1}{4} \sum_r \chi_{rr} + \frac{1}{4} \sum_{r>s} \chi_{rs}.$$

$$\Delta^{fund} = G(0) - ZPVE^{fund} = -\frac{3}{4} \sum_r \chi_{rr} - \frac{1}{4} \sum_{r>s} \chi_{rs}.$$

## Characterization of Stationary Points

- A *stationary point* is a geometry  $\tilde{\mathbf{q}}$  for which the gradient  $\partial E_e(\tilde{\mathbf{q}})/\partial \tilde{q}_i$  for all coordinates  $\tilde{q}_i$ : can be a (global or local) PES minimum, transition state, or higher order saddle point
- The *Hessian Index* is the number of negative force constants (corresponding to imaginary vibrational frequencies, often printed as negative frequencies)
- For a minimum, verify that there are no imaginary frequencies
- For a transition state, verify there is exactly one unique imaginary frequency

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*****
Running Q-Chem

cicero.chemistry.gatech.edu on Fri Dec 15 14:51:29 EST 2000.

Version: /usr/local/qchem.6.3.00/bin/qchem
        /usr/local/qchem.6.3.00/exe/progman.exe

*****
                Welcome to Q-Chem
        A Quantum Leap Into The Future Of Chemistry

J. Kong, C. A. White, A. I. Krylov, C. D. Sherrill,
R. D. Adamson, T. R. Furlani, M. S. Lee, A. M. Lee,
S. R. Gwaltney, T. R. Adams, H. Dachsel, W. M. Zhang,
P. P. Korambath, C. Ochsenfeld, A. T. B. Gilbert,
G. S. Kedziora, D. R. Maurice, N. Nair, Y. Shao,
N. A. Besley, P. E. Maslen, J. P. Dombroski, J. Baker,
E. F. C. Byrd, T. Van Voorhis, M. Oumi, S. Hirata,
C. P. Hsu, N. Ishikawa, J. Florian, A. Warshel,
B. G. Johnson, P. M. W. Gill, M. Head-Gordon, J. A. Pople,
Q-Chem, Version 2.0, Q-Chem, Inc., Export, PA (2000).

Intel x86 Linux Version

=====
User input:
=====
$comment
Water
$end

$molecule
0 1
O
H1 O OH
H2 O OH H1 HOH

OH = 0.989276
HOH = 100.0198
$end

$rem
JOBTYPE          FREQ
EXCHANGE         HF
CORRELATION      NONE
BASIS            STO-3G
SCF_FINAL_PRINT  2          Energies plus MOS
$end

=====
Processing $rem in the input.

#####
# Entering fldman.exe on Fri Dec 15 14:51:30 2000 #
#####

-----
Standard Nuclear Orientation (Angstroms)
-----
I      Atom      X      Y      Z
-----
1      O      0.000000  0.000000  0.127153
2      H      -0.757939  0.000000  -0.508611
3      H      0.757939  0.000000  -0.508611
-----

Molecular Point Group          C2v  NOp = 4
Largest Abelian Subgroup      C2v  NOp = 4
Nuclear Repulsion Energy =    8.9077080996 hartrees

```

```

There are      5 alpha and      5 beta electrons
Requested basis set is STO-3G
There are 4 shells and 7 basis functions
A cutoff of 1.0D-10 yielded      10 shell pairs
There are      34 function pairs

#####
# Entering gesman.exe on Fri Dec 15 14:51:30 2000 #
#####

Smallest overlap matrix eigenvalue = 3.63E-01
Multipole matrices computed through 2nd order
Guess from superposition of atomic densities
Warning: Energy on first SCF cycle will be non-variational

#####
# Entering scfman.exe on Fri Dec 15 14:51:30 2000 #
#####

A restricted Hartree-Fock SCF calculation will be
performed using Pulay DIIS extrapolation
SCF converges when DIIS error is below 1.0E-08

-----
Cycle      Energy      DIIS Error
-----
1      -74.6061773729      4.01E-01  000000
2      -74.9233301871      6.78E-02  000000
3      -74.9653425178      8.77E-03  000000
4      -74.9658786194      1.57E-03  000000
5      -74.9659011332      2.76E-05  000000
6      -74.9659011457      1.08E-05  000000
7      -74.9659011480      7.89E-08  000000
8      -74.9659011480      2.72E-08  000000
9      -74.9659011480      1.03E-08  000000
10     -74.9659011480      3.60E-10  000000  Convergence criterion met

-----
SCF time: CPU 0.08 s wall 0.00 s

Final Alpha MO Eigenvalues
      1      2      3      4      5      6
1 -20.2515674 -1.2576205 -0.5939092 -0.4597630 -0.3926280 0.5819312
7
1 0.6928062
Final Alpha MO Coefficients
      1      2      3      4      5      6
1 0.9942161 -0.2337592 0.0000000 -0.1040490 0.0000000 0.1258359
2 0.0258496 0.8443968 -0.0000000 0.5382402 -0.0000000 -0.8203442
3 0.0000000 0.0000000 0.6127098 0.0000000 -0.0000000 0.0000000
4 -0.0000000 0.0000000 0.0000000 0.0000000 1.0000000 -0.0000000
5 -0.0041658 -0.1228911 -0.0000000 0.7558818 -0.0000000 0.7635987
6 -0.0055852 0.1556051 -0.4492190 -0.2950597 0.0000000 0.7692443
7 -0.0055852 0.1556051 0.4492190 -0.2950597 0.0000000 0.7692443
7
1 -0.0000000
2 0.0000000
3 0.9598513
4 -0.0000000
5 -0.0000000
6 0.8147663
7 -0.8147663
Final Alpha density matrix.
      1      2      3      4      5      6
1 1.0539352 -0.2276888 0.0000000 -0.0000000 -0.0540635 -0.0112263
2 -0.2276888 1.0033767 -0.0000000 0.0000000 0.3029695 -0.0275650
3 0.0000000 -0.0000000 0.3754133 -0.0000000 -0.0000000 -0.2752409
4 -0.0000000 0.0000000 -0.0000000 1.0000000 -0.0000000 -0.0000000
5 -0.0540635 0.3029695 -0.0000000 -0.0000000 0.5864768 -0.2421295
6 -0.0112263 -0.0275650 -0.2752409 -0.0000000 -0.2421295 0.3131021
7 -0.0112263 -0.0275650 0.2752409 -0.0000000 -0.2421295 -0.0904933
7

```

```

1 -0.0112263
2 -0.0275650
3 0.2752409
4 -0.0000000
5 -0.2421295
6 -0.0904933
7 0.3131021

#####
# Entering anlman.exe on Fri Dec 15 14:51:30 2000 #
#####

Analysis of SCF Wavefunction

-----
Orbital Energies (a.u.) and Symmetries
-----

Alpha MOs, Restricted
-- Occupied --
-20.252 -1.258 -0.594 -0.460 -0.393
 1 A1  2 A1  1 B1  3 A1  1 B2
-- Virtual --
 0.582 0.693
 4 A1  2 B1

Beta MOs, Restricted
-- Occupied --
-20.252 -1.258 -0.594 -0.460 -0.393
 1 A1  2 A1  1 B1  3 A1  1 B2
-- Virtual --
 0.582 0.693
 4 A1  2 B1

-----
Mulliken Net Atomic Charges

Atom          Charge (a.u.)
-----
 1 O           -0.330636
 2 H            0.165318
 3 H            0.165318

Sum of atomic charges = 0.000000

-----
Cartesian Multipole Moments

Charge (ESU x 10^10)
0.0000
Dipole Moment (Debye)
  X 0.0000  Y 0.0000  Z -1.7094
Tot 1.7094
Quadrupole Moments (Debye-Ang)
  XX -4.4859  XY 0.0000  YY -6.1256
  XZ 0.0000  YZ 0.0000  ZZ -5.3331
Octapole Moments (Debye-Ang^2)
  XXX 0.0000  XXY 0.0000  XYY 0.0000
  YYY 0.0000  XXZ -0.5313  XYZ 0.0000
  YYZ 0.0191  XZZ 0.0000  YZZ 0.0000
  ZZZ -0.1747
Hexadecapole Moments (Debye-Ang^3)
  XXXX -6.7321  XXXY 0.0000  XYYX -1.8081
  YYYY 0.0000  YYYY -3.2652  XXXZ 0.0000
  XXYZ 0.0000  XYYZ 0.0000  YYYZ 0.0000
  XXZZ -1.7388  XYZZ 0.0000  YZZZ -1.4596
  XZZZ 0.0000  YZZZ 0.0000  ZZZZ -5.2191

#####

```

```

# Entering drvman.exe on Fri Dec 15 14:51:30 2000 #
#####

Calculating MO derivatives via CPHF
 1 10 2 0.009892 0.007094
 2 12 0 0.000000 0.000000 Roots Converged

Calculating analytic Hessian of the SCF energy
Polarizability Matrix (a.u.)
 1 2 3
 1 -5.5054487 -0.0000000 -0.0000000
 2 -0.0000000 -0.0400454 -0.0000000
 3 -0.0000000 -0.0000000 -2.5654448

Direct stationary perturbation theory relativistic correction:

rels = 0.031221304489
relv = -0.096844180832
rel2e = 0.023322614724
E_rel = -0.042300261619

Hessian of the SCF Energy
 1 2 3 4 5 6
 1 0.8044023 0.0000000 -0.0000000 -0.4022012 -0.0000000 -0.3374251
 2 0.0000000 -0.001352 0.0000000 0.0000000 0.0000676 -0.0000000
 3 -0.0000000 0.0000000 0.6352335 -0.2165519 -0.0000000 -0.3176168
 4 -0.4022012 -0.0000000 -0.2165519 0.4391751 0.0000000 0.2769885
 5 -0.0000000 0.0000676 -0.0000000 0.0000000 -0.0000838 0.0000000
 6 -0.3374251 -0.0000000 -0.3176168 0.2769885 0.0000000 0.3002754
 7 -0.4022012 -0.0000000 0.2165519 -0.0369740 -0.0000000 0.0604366
 8 0.0000000 0.0000676 -0.0000000 -0.0000000 0.0000162 0.0000000
 9 0.3374251 -0.0000000 -0.3176168 -0.0604366 -0.0000000 0.0173414
 7 8 9
 1 -0.4022012 0.0000000 0.3374251
 2 -0.0000000 0.0000676 -0.0000000
 3 0.2165519 -0.0000000 -0.3176168
 4 -0.0369740 -0.0000000 -0.0604366
 5 -0.0000000 0.0000162 -0.0000000
 6 0.0604366 0.0000000 0.0173414
 7 0.4391751 0.0000000 -0.2769885
 8 0.0000000 -0.0000838 0.0000000
 9 -0.2769885 0.0000000 0.3002754
Gradient time: CPU 0.59 s wall 1.00 s

#####
# Entering vibman.exe on Fri Dec 15 14:51:31 2000 #
#####

*****
**
** VIBRATIONAL ANALYSIS
**
** -----
** VIBRATIONAL FREQUENCIES (CM**-1) AND NORMAL MODES
** INFRARED INTENSITIES (KM/MOL)
**
*****

Frequency: 2169.95 4141.60 4392.63
IR Active: YES YES YES
IR Intens: 7.245 44.303 29.972
Raman Active: YES YES YES

X Y Z X Y Z X Y Z
O 0.000 0.000 -0.069 0.000 0.000 0.052 -0.068 0.000 0.000
H -0.448 0.000 0.545 -0.570 0.000 -0.416 0.540 0.000 0.453
H 0.448 0.000 0.545 0.570 0.000 -0.416 0.540 0.000 -0.453

STANDARD THERMODYNAMIC QUANTITIES AT 298.18 K AND 1.00 ATM

This Molecule has 0 Imaginary Frequencies

```